

Orbit determination with the two-body integrals

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- The orbit determination problem
- The integrals of Kepler's problem
- Optical observations and attributables
- Algebraic methods for preliminary orbits
- A test case
- Conclusions and open problems

The orbit determination problem: old and new



- **observations**: single observations per night
- **preliminary orbits**: Gauss or Laplace methods
- **least squares fit**: refinement of the preliminary orbits

- **observations**: very short arcs (VSAs)
- **identification of VSAs for preliminary orbits**:
new methods needed
- **least squares fit**: still necessary

Kepler's problem and its integrals

$$\ddot{\mathbf{r}} = -\mu \frac{\mathbf{r}}{|\mathbf{r}|^3}, \quad \mathbf{r} \in \mathbb{R}^3, \quad \mu > 0.$$

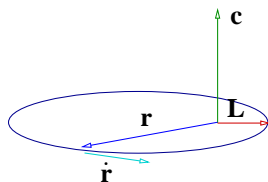
Kepler's problem

INTEGRALS OF MOTIONS

$$\mathbf{c} = \mathbf{r} \times \dot{\mathbf{r}} \quad \text{angular momentum}$$

$$\mathcal{E} = \frac{1}{2} |\dot{\mathbf{r}}|^2 - \frac{\mu}{|\mathbf{r}|} \quad \text{energy}$$

$$\mathbf{L} = \frac{1}{\mu} \dot{\mathbf{r}} \times \mathbf{c} - \frac{\mathbf{r}}{|\mathbf{r}|} \quad \text{Laplace-Lenz vector}$$



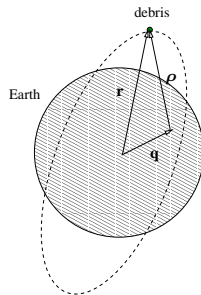
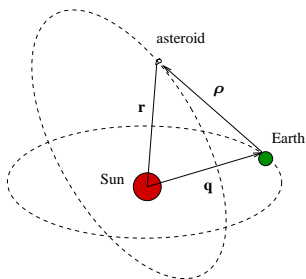
Only 5 scalar integrals are independent:

$$\mathbf{L} \cdot \mathbf{c} = 0$$

$$2|\mathbf{c}|^2 \mathcal{E} + \mu^2(1 - |\mathbf{L}|^2) = 0$$

Geometry of the observations

Asteroid and space debris observations



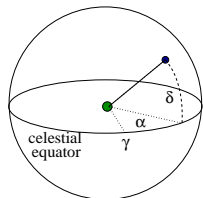
$$\mathbf{r} = \boldsymbol{\rho} + \mathbf{q}$$

$$\boldsymbol{\rho} = \rho \hat{\mathbf{e}}^\rho$$

$$\hat{\mathbf{e}}^\rho = \text{observation direction}$$

To compute an orbit we have to determine $\mathbf{r}, \dot{\mathbf{r}}$ at some time.

Optical observations of a Solar system body



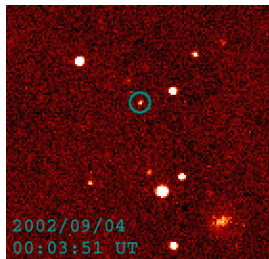
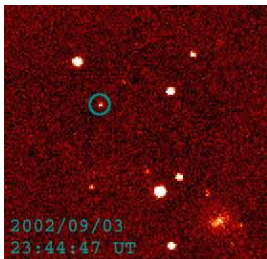
α =right ascension

δ = declination

An optical observation consists of two angles (α, δ) on the celestial sphere.

The topocentric distance ρ of the the observed body is unknown.

DETECTION OF AN ASTEROID



Given a short arc of optical observations

$$(\alpha_i, \delta_i) \in S^1 \times (-\pi/2, \pi/2) \text{ at times } t_i, i = 1 \dots m,$$

of a Solar system body we can compute an **attributable**

$$\mathcal{A} = (\alpha, \delta, \dot{\alpha}, \dot{\delta})$$

at time $\bar{t} = \frac{1}{m} \sum_i t_i$ by linear or quadratic interpolation.

The radial distance ρ and the radial velocity $\dot{\rho}$ are completely undetermined.

Algebraic methods for the linkage problem

Linkage/correlation: identification of two short arcs as related to the same observed body. Using the data of both arcs we can try to compute an orbit.

Our **main idea** is to perform linkage by solving polynomial equations using the first integrals of the Kepler problem. This allows a global control on the solutions.

Taff and Hall (1977): first suggested to use the integrals c, \mathcal{E} to write equations for the linkage problem.

Gronchi, Dimare and Milani (2010): proposed to solve the equations by Taff and Hall by algebraic methods.

Gronchi, Farnocchia and Dimare (2011): proposed to solve different equations, using a suitable projection of the Laplace-Lenz vector in place of the energy.

Keplerian integrals as function of $(\rho, \dot{\rho})$

ANGULAR MOMENTUM:

$$\mathbf{c}(\rho, \dot{\rho}) = \mathbf{r} \times \dot{\mathbf{r}} = \mathbf{D}\dot{\rho} + \mathbf{E}\rho^2 + \mathbf{F}\rho + \mathbf{G}$$

with

$$\begin{aligned}\mathbf{D} &= \mathbf{q} \times \hat{\mathbf{e}}^\rho \\ \mathbf{E} &= \dot{\alpha} \cos \delta \hat{\mathbf{e}}^\rho \times \hat{\mathbf{e}}^\alpha + \dot{\delta} \hat{\mathbf{e}}^\rho \times \hat{\mathbf{e}}^\delta \\ \mathbf{F} &= \dot{\alpha} \cos \delta \mathbf{q} \times \hat{\mathbf{e}}^\alpha + \dot{\delta} \mathbf{q} \times \hat{\mathbf{e}}^\delta + \hat{\mathbf{e}}^\rho \times \dot{\mathbf{q}} \\ \mathbf{G} &= \mathbf{q} \times \dot{\mathbf{q}}\end{aligned}$$

where $\hat{\mathbf{e}}^\rho, \hat{\mathbf{e}}^\alpha, \hat{\mathbf{e}}^\delta$ are unit vectors depending only on α, δ .
Thus $\mathbf{D}, \mathbf{E}, \mathbf{F}, \mathbf{G}$ depend only on the attributable \mathcal{A} and on the motion of the observer $\mathbf{q}, \dot{\mathbf{q}}$ at time \bar{t} .

Keplerian integrals as function of $(\rho, \dot{\rho})$

ENERGY:

$$\mathcal{E} = \frac{1}{2}|\dot{\mathbf{r}}|^2 - \frac{\mu}{|\mathbf{r}|}$$

LAPLACE-LENZ VECTOR:

$$\mu\mathbf{L}(\rho, \dot{\rho}) = \dot{\mathbf{r}} \times \mathbf{c} - \mu \frac{\mathbf{r}}{|\mathbf{r}|} = \left(|\dot{\mathbf{r}}|^2 - \frac{\mu}{|\mathbf{r}|} \right) \mathbf{r} - (\dot{\mathbf{r}} \cdot \mathbf{r}) \dot{\mathbf{r}}$$

where

$$|\mathbf{r}| = (\rho^2 + |\mathbf{q}|^2 + 2\rho\mathbf{q} \cdot \hat{\mathbf{e}}^\rho)^{1/2},$$

$$|\dot{\mathbf{r}}|^2 = \dot{\rho}^2 + (\dot{\alpha}^2 \cos^2 \delta + \dot{\delta}^2)\rho^2 + 2\dot{\mathbf{q}} \cdot \hat{\mathbf{e}}^\rho \dot{\rho} + 2\dot{\mathbf{q}} \cdot (\dot{\alpha} \cos \delta \hat{\mathbf{e}}^\alpha + \dot{\delta} \hat{\mathbf{e}}^\delta)\rho + |\dot{\mathbf{q}}|^2$$

$$\dot{\mathbf{r}} \cdot \mathbf{r} = \rho \dot{\rho} + \mathbf{q} \cdot \hat{\mathbf{e}}^\rho \dot{\rho} + (\dot{\mathbf{q}} \cdot \hat{\mathbf{e}}^\rho + \mathbf{q} \cdot \hat{\mathbf{e}}^\alpha \dot{\alpha} \cos \delta + \mathbf{q} \cdot \hat{\mathbf{e}}^\delta \dot{\delta})\rho + \dot{\mathbf{q}} \cdot \mathbf{q}$$

Equating the angular momentum

Given two attributables $\mathcal{A}_1, \mathcal{A}_2$ at times \bar{t}_1, \bar{t}_2 , equating the angular momentum at the two times we obtain

$$\mathbf{D}_1 \dot{\rho}_1 - \mathbf{D}_2 \dot{\rho}_2 = \mathbf{J}(\rho_1, \rho_2) \quad (1)$$

where

$$\mathbf{J}(\rho_1, \rho_2) = \mathbf{E}_2 \rho_2^2 - \mathbf{E}_1 \rho_1^2 + \mathbf{F}_2 \rho_2 - \mathbf{F}_1 \rho_1 + \mathbf{G}_2 - \mathbf{G}_1 .$$

By scalar multiplication of (1) with $\mathbf{D}_1 \times \mathbf{D}_2$ we perform elimination of the variables $\dot{\rho}_1, \dot{\rho}_2$; this yields

$$q(\rho_1, \rho_2) \stackrel{\text{def}}{=} \mathbf{D}_1 \times \mathbf{D}_2 \cdot \mathbf{J}(\rho_1, \rho_2) = 0$$

Equating the energy

By vector multiplication of (1) with \mathbf{D}_1 and \mathbf{D}_2 , projecting on $\mathbf{D}_1 \times \mathbf{D}_2$ we obtain

$$\dot{\rho}_1(\rho_1, \rho_2) = \frac{(\mathbf{J} \times \mathbf{D}_2) \cdot (\mathbf{D}_1 \times \mathbf{D}_2)}{|\mathbf{D}_1 \times \mathbf{D}_2|^2}; \quad \dot{\rho}_2(\rho_1, \rho_2) = \frac{(\mathbf{J} \times \mathbf{D}_1) \cdot (\mathbf{D}_1 \times \mathbf{D}_2)}{|\mathbf{D}_1 \times \mathbf{D}_2|^2}$$

Substituting into $\mathcal{E}_1 = \mathcal{E}_2$, rearranging the terms and squaring twice we obtain a polynomial equation

$$p(\rho_1, \rho_2) = 0$$

with total degree 24.

Intersections between the curves

We study the semi-algebraic problem

$$\begin{cases} p(\rho_1, \rho_2) = 0 \\ q(\rho_1, \rho_2) = 0 \end{cases}, \quad \rho_1, \rho_2 > 0 \quad (2)$$

by **resultant theory**.

We can write

$$p(\rho_1, \rho_2) = \sum_{j=0}^{20} a_j(\rho_2) \rho_1^j,$$
$$q(\rho_1, \rho_2) = b_2 \rho_1^2 + b_1 \rho_1 + b_0(\rho_2)$$

for some coefficients a_i, b_j , depending only on ρ_2 .

Consider the resultant $Res(\rho_2)$ of p, q with respect to ρ_1 : it is a **48 degree polynomial** defined as the determinant of the Sylvester matrix

$$\begin{pmatrix} a_{20} & 0 & b_2 & 0 & \dots & \dots & 0 \\ a_{19} & a_{20} & b_1 & b_2 & 0 & \dots & 0 \\ \vdots & \vdots & b_0 & b_1 & b_2 & \dots & \vdots \\ \vdots & \vdots & 0 & b_0 & b_1 & \dots & \vdots \\ a_0 & a_1 & \vdots & \vdots & \vdots & b_0 & b_1 \\ 0 & a_0 & 0 & 0 & 0 & 0 & b_0 \end{pmatrix}.$$

The ρ_2 component of a solution of (2) must be a root of $Res(\rho_2)$.

Computation of the solutions

- Compute the positive roots ρ_2 of $Res(\rho_2)$.
- For each root find the corresponding values of $\rho_1, \dot{\rho}_1, \dot{\rho}_2$.
- Discard spurious solutions, obtained by squaring.
- Compute the related Keplerian orbits at times

$$\tilde{t}_i = \bar{t}_i - \frac{\rho_i}{c}, i = 1, 2$$

corrected by aberration.

Selecting the solutions I: compatibility conditions

The knowledge of the angular momentum vector and of the energy at a given time yields the values of

$$a, e, I, \Omega .$$

From $\mathbf{c}_1 = \mathbf{c}_2$, $\mathcal{E}_1 = \mathcal{E}_2$ we obtain the same values of a, e, I, Ω at times \tilde{t}_1, \tilde{t}_2 , but we must check for the *compatibility conditions*

$$\omega_1 = \omega_2, \quad \ell_1 = \ell_2 + n(\tilde{t}_1 - \tilde{t}_2), \quad (3)$$

where n is the *mean motion*.

Problem: how to take into account the errors in the observations?

Selecting the solutions I: compatibility conditions

Consider the map

$$(\mathcal{A}_1, \mathcal{A}_2) \xrightarrow{\Psi} (\mathcal{A}_1, \rho_1, \dot{\rho}_1, \Delta_{1,2}), \quad \Delta_{1,2} = (\Delta\omega, \Delta\ell)$$

giving **the orbit at time \tilde{t}_1** , and **the difference in ω, ℓ** .

Check whether $\Delta_{1,2} = 0$ is compatible with the observational errors by covariance propagation through the map Ψ .

This algorithm also allows to **define covariance matrices for the preliminary orbits that we compute**.

Equating a component of the Laplace-Lenz vector

It is possible to reduce the algebraic degree of the linkage problem by writing different equations (i.e. using different integrals). In [Gronchi, Farnocchia, Dimare \(2011\)](#) we select **a suitable component of the Laplace-Lenz vector** in place of the energy.

Given $\mathcal{A}_1, \mathcal{A}_2$ we equate $\mathbf{L}_1, \mathbf{L}_2$ projected along $\mathbf{v} = \hat{\mathbf{e}}_2^{\rho} \times \mathbf{q}_2$:

$$\mathbf{L}_1(\rho_1, \dot{\rho}_1) \cdot \mathbf{v} = \mathbf{L}_2(\rho_2, \dot{\rho}_2) \cdot \mathbf{v} .$$

Equating a component of the Laplace-Lenz vector

We have

$$\left(|\dot{\mathbf{r}}_1|^2 - \frac{\mu}{|\mathbf{r}_1|} \right) (\mathbf{r}_1 \cdot \mathbf{v}) - (\dot{\mathbf{r}}_1 \cdot \mathbf{r}_1)(\dot{\mathbf{r}}_1 \cdot \mathbf{v}) = -(\dot{\mathbf{r}}_2 \cdot \mathbf{r}_2)(\dot{\mathbf{r}}_2 \cdot \mathbf{v}) .$$

Rearranging the terms and squaring we obtain

$$\tilde{p}(\rho_1, \rho_2) \stackrel{\text{def}}{=} \mu^2 (\mathbf{r}_1 \cdot \mathbf{v})^2 - |\mathbf{r}_1|^2 \left\{ [|\dot{\mathbf{r}}_1|^2 \mathbf{r}_1 - (\dot{\mathbf{r}}_1 \cdot \mathbf{r}_1) \dot{\mathbf{r}}_1 + (\dot{\mathbf{r}}_2 \cdot \mathbf{r}_2) \dot{\mathbf{r}}_2] \cdot \mathbf{v} \right\}^2 = 0 .$$

\tilde{p} is a polynomial of **total degree 10** in ρ_1, ρ_2 , therefore

$$\tilde{p}(\rho_1, \rho_2) = 0, \quad q(\rho_1, \rho_2) = 0, \quad \rho_1, \rho_2 > 0 .$$

has **degree 20**.

Selecting the solutions II: attribution

We can select the solutions of the linkage problem also by means of the **attribution** algorithm.

Let \mathcal{E}_1 be a set of orbital elements at time t_1 , with covariance matrix Γ_1 . Propagate orbit and covariance to the epoch \bar{t}_2 of \mathcal{A}_2 , with covariance $\Gamma_{\mathcal{A}_2}$. Then extract a **predicted attributable** \mathcal{A}_p , at time \bar{t}_2 , with covariance $\Gamma_{\mathcal{A}_p}$.

We can compare $\mathcal{A}_p, \Gamma_{\mathcal{A}_p}$ with $\mathcal{A}_2, \Gamma_{\mathcal{A}_2}$ by defining an **identification penalty** χ_4 , that gives a measure of the '*price to pay*' assuming the observations of both attributables belong to the same celestial body.

A test case: (99942) Apophis

DATA: two sets of observations with mean epochs

$$\bar{t}_1 = 53175.59, \bar{t}_2 = 53357.45 \text{ MJD.}$$

SOLUTIONS:

	ρ_1	ρ_2
1	0.78987	0.04345
2	1.13777	0.09569

Attribution: $\chi_4(1) = 3230925.94$, $\chi_4(2) = 2.29$. We select the second one:

$$a = 0.923, \quad e = 0.189, \quad I = 3.287$$
$$\Omega = 204.912, \quad \omega = 124.778, \quad \ell = 249.003$$

at epoch $t_1 = 53175.59$. We compare it with the known orbit:

$$a = 0.922, \quad e = 0.191, \quad I = 3.333$$
$$\Omega = 204.575, \quad \omega = 126.176, \quad \ell = 247.500$$

Laplace-Lenz versus Energy

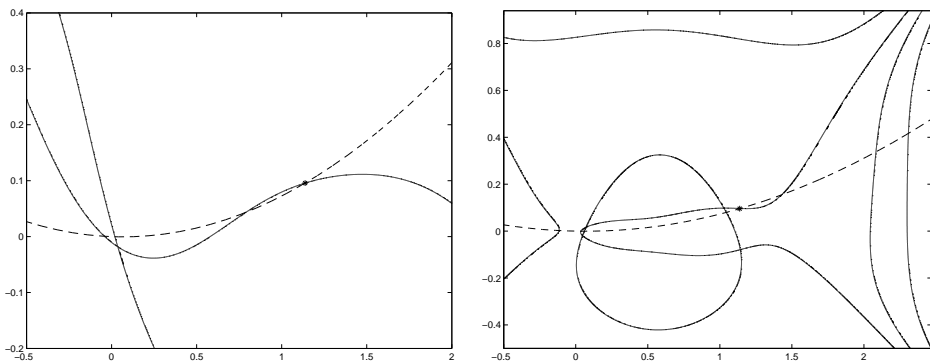


Figure: Test case with (99942) Apophis. Left: with Laplace-Lenz vector. Right: with energy.

In [Dimare et al. 2011](#) a comparison of the two methods is shown for a LEO population.

Conclusions and open problems

The investigation of new orbit determination methods is essential to deal with modern observations. Algebraic methods are suitable and effective for the linkage problem.

J_2 EFFECTS: in case of space debris the effect of oblateness of the Earth cannot be neglected: in Farnocchia et al. (2010) a solution is proposed.

RADAR OBSERVATIONS: We can define a radar attributable $\mathcal{A}_{rad} = (\alpha, \delta, \rho, \dot{\rho})$ at time \bar{t} . Here $\dot{\alpha}, \dot{\delta}$ remain completely undetermined.

Problem: large uncertainties of the angles α, δ . Can we use this data type?

If yes then we have an explicit solution of the linkage problem for the mixed case of 1 optical and 1 radar attributable.

- [1] **L. Dimare et al.** *Innovative system of very wide field optical sensors for space surveillance in the LEO region*, AMOS conf. 2011
- [2] **D. Farnocchia et al.:** *Innovative methods of correlation and orbit determination for space debris*, CMDA, 2010
- [3] **G. F. Gronchi, L., Dimare, A., Milani:** *Orbit determination with the two-body integrals*, CMDA, 2010
- [4] **G. F. Gronchi, D. Farnocchia, L., Dimare:** *Orbit determination with the two-body integrals. II*, CMDA, 2011
- [5] **L. G. Taff, D. L. Hall:** *The use of angles and angular rates. I - Initial orbit determination*, CMDA, 1977